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Defect-related carrier transport in CdTe-based compounds: comparison with hybrid perovskites

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- 3) Savannah River Natl Lab, Aiken, SC 29808 USA
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Thanks:

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- Savannah River National Laboratory R.B. James, U. N. Roy
- Freiburger Materialforschungszentrum (FMF) M. Fiederle, A. Fauler, V. Babentsov
- University of Surrey P. Sellin
- IMEM, Parma A. Zappettini
- University of Bologna B. Fraboni, A. Cavallini
- Korea University K.H.Kim
- Redlen technologies G. Prekas, K. Iniewski

History of CdTe crystal growth and defect studies at the Institute of Physics, Charles University

1962 – 1st paper (Growing of CdTe single crystals by static sublimation.., P.Höschl, C.Konak, Czech.J.Phys.)

1977 – paper P. Höschl et al., Rev. Phys. Appl.



1985 –2002 HgCdTe bulk crystals for infrared detectors
CdZnTe crystals as substrates for HgCdTe MBE
epitaxial layers

CdTe
CdZnTe
CdZnTeSe

For X-ray and gamma-ray detectors

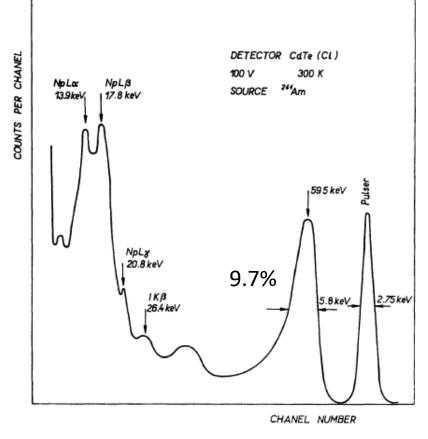


Fig. 5. — ²⁴¹Am gamma-ray spectrum at room temperature.

Defects and their impact on detector performance

Methods

photoluminescence

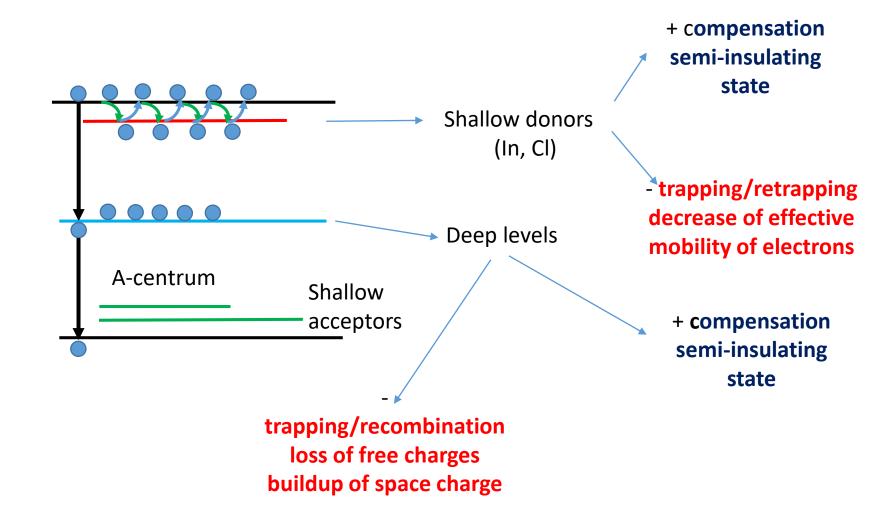
TEES, PICTS

photoconductivity

Pockels effect

Transient current technique (L-TCT)

Photo-Hall effect spectroscopy



Achievement of the semi-insulating state

 $N_D \sim 10^{15} cm^{-3}$

 $N_{DEEP} \sim 10^{11} cm^{-3}$ $\sigma \sim 10^{-13} - 10^{-14} cm^{2}$

 $N_A \sim 10^{15} cm^{-3}$

	Fermi energy	→	$N_{DEEP} > N_D - N_A cm^{-3}$



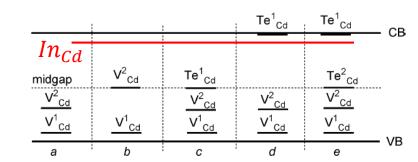
Model

Slow cooling at Te-rich conditions

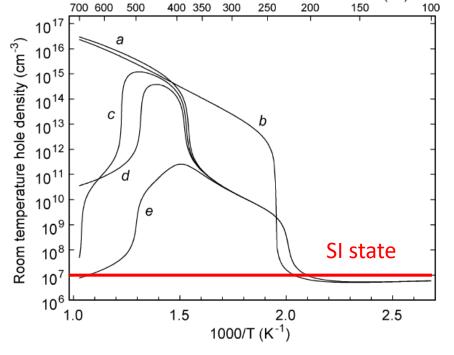
Reduction of point defects by precipitation



Semi-insulating state is achieved at ~200K



R. Grill, J. Franc, P. Höschl, I. Turkevych, E.Belas, P. Moravec, IEEE Trans. Nucl. Sci. 52, 1925, 2005.



T (°C)

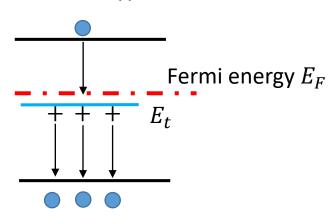
 V_{Cd} , $Te_{Cd} \rightarrow Te_{ppt}$

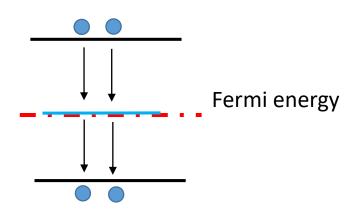
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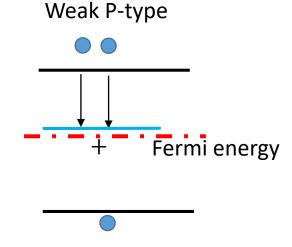
Midgap level

$$\sigma_{t,e} = \frac{1}{\sigma_e v_{th} N_t f} \qquad f = \frac{1}{1 + e^{E_t - R}}$$

Weak N-type







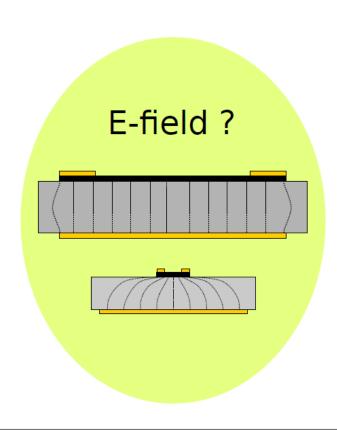
Weak electron trapping, Good electron mobility

Strong hole trapping, Buildup of positive space charge Recombination center

Due to a higher electron than hole mobility trapping of electrons is weaker than the trapping of holes Moderate electron trapping Weak hole trapping

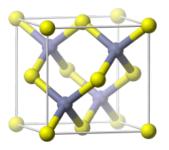
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Electric field measurements – Pockels effect



Gauss law: E-field \iff space charge

CdTe, CdZnTe, CdZnTeSe . . . optically isotropic crystals of 43m symmetry (FCC) showing Pockels effect (E-field induced birefringence)



E-field	Crystal	
0	isotropic	
≠ 0	anisotropic	

Taylor:

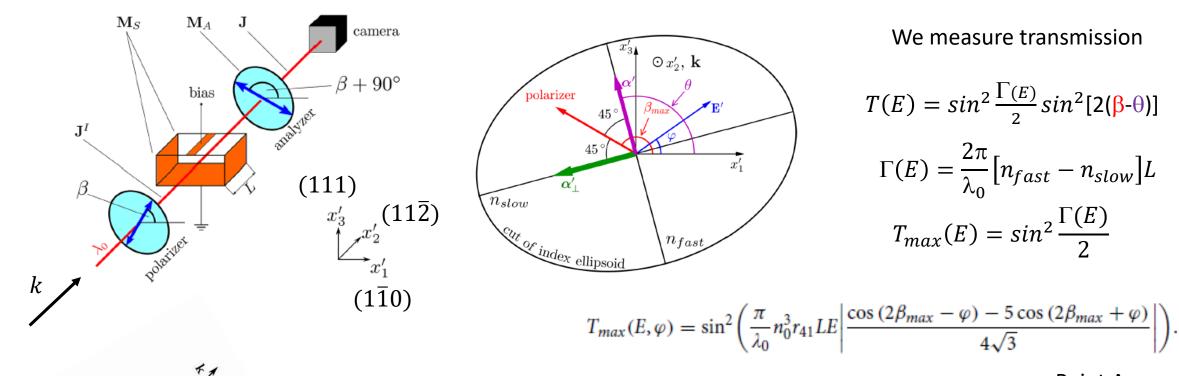
$$n(E) = n + a_1 E + \frac{1}{2} a_2 E^2 + \dots$$

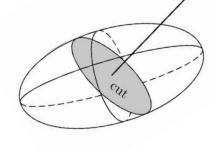
with electrooptic coefficients r and s:

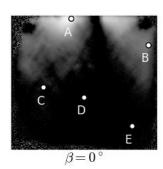
$$(n(E) \approx n - \frac{1}{2}rn^3E - \frac{1}{2}sn^3E^2$$

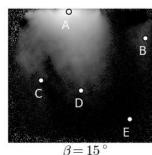
Pockels effect Kerr effect (linear) (quadratic)

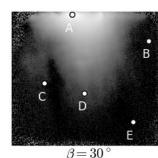
Pockels effect - arbitrary direction of E-field

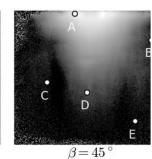


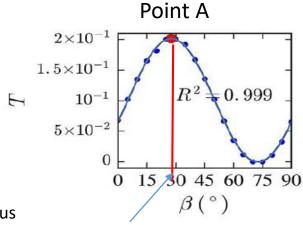








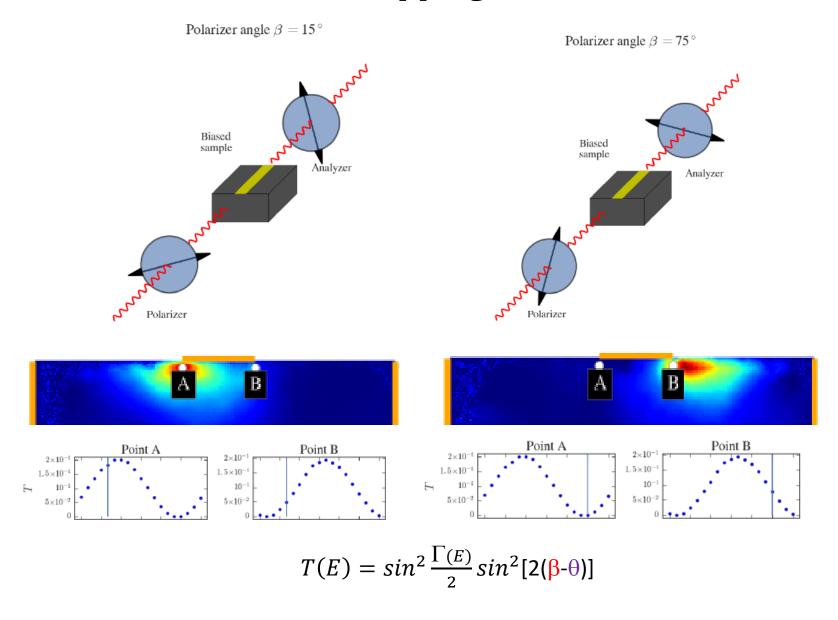




 β_{max}

Václav Dědič, Tomáš Fridrišek, Jan Franc, Jan Kunc, Martin Rejhon, Utpal N. Roy, Ralph B. James, Mapping of inhomogenous quasi-3D electrostatic field in electro-optic materials, Scientific Reports (2021) 11:2154

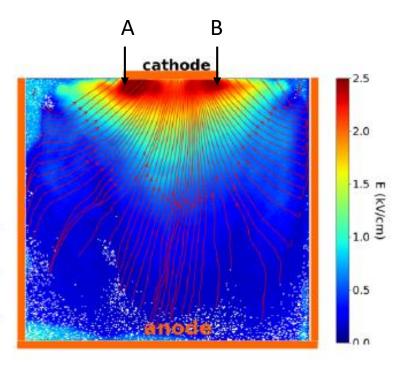
Mapping of electric fields



Evaluation of the E-field map:

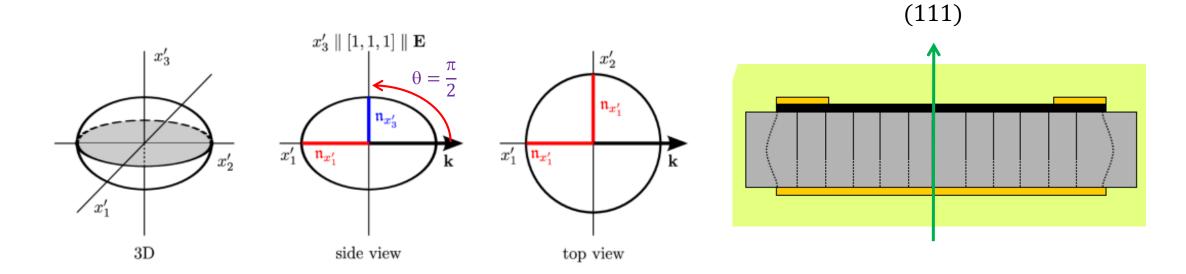
- 1. Each pixel must be treated as a separate sample
- Maximum of transmission T is found rotating the crossed polarizer-analyzer set
- 3. E-field calculated from

$$T_{max}(E) = \sin^2 \frac{\Gamma(E)}{2}$$



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Pockels effect – E-field is parallel to (111)

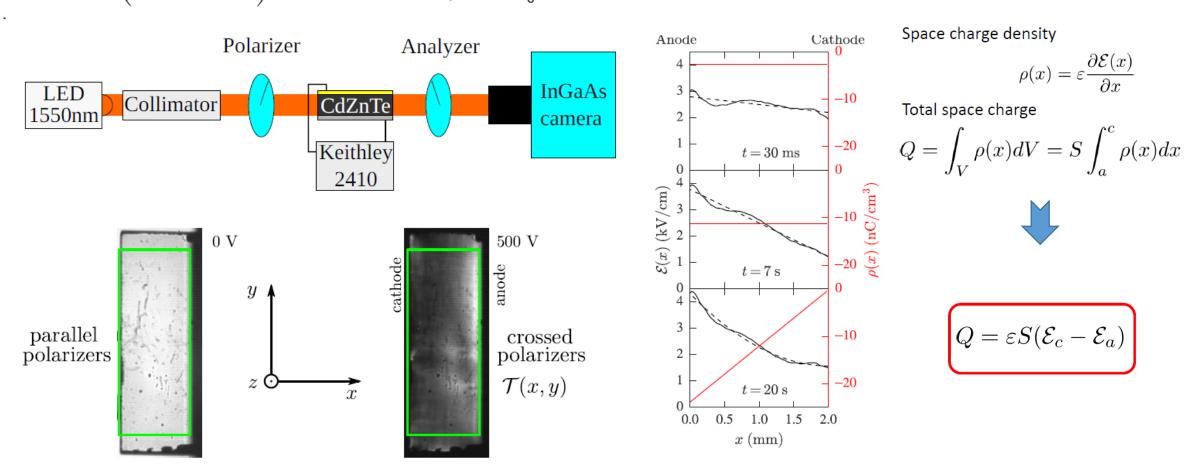


$$\theta = \frac{\pi}{2}, \text{ for } \beta = \frac{\pi}{4} \implies 2(\beta - \theta) = -\frac{\pi}{2} \implies \sin^2[2(\beta - \theta)] = 1 \implies T(E) = T_{max}(E) = \sin^2\frac{\Gamma(E)}{2}$$

Simple situation, with polarizer angle $\beta = \frac{\pi}{4}$, analyzer angle $\beta = \frac{3\pi}{4}$, T_{max} is measured

Electric field evaluation

$$T = \sin^2\left(\frac{\sqrt{3}}{2} \frac{\pi r_{41} n_0^3 EL}{\lambda_0}\right) \Rightarrow E(x_1', x_3') = \frac{2}{\sqrt{3}} \frac{\lambda_0}{\pi r_{41} n_0^3 L} \arcsin\sqrt{T(x_1', x_3')} = \alpha_P \arcsin\sqrt{T(x_1', x_3')}$$



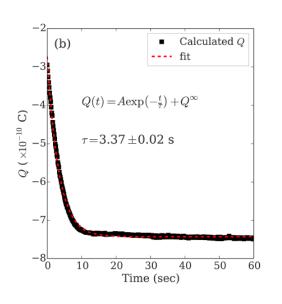
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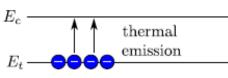
Experimental methods based on Pockels effect I

E-field DLTS

(DLTS = deep level transient spectroscopy)

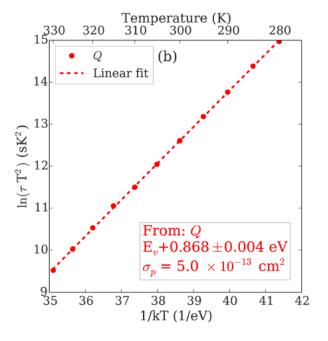
- Subject of study: thermal emission of carriers from traps
- Temporal and temperature dependence of the E-field (reflecting space charge) after switchin ON the bias or switching OFF the illumination (light assisted trap filling)
- Output: activation energy and capture cross-sections of traps (Arrhenius plot)⇒deep levels responsible for space charge accumulation
- Advantage: higher sensitivity than current techniques, direct space charge sign estimation





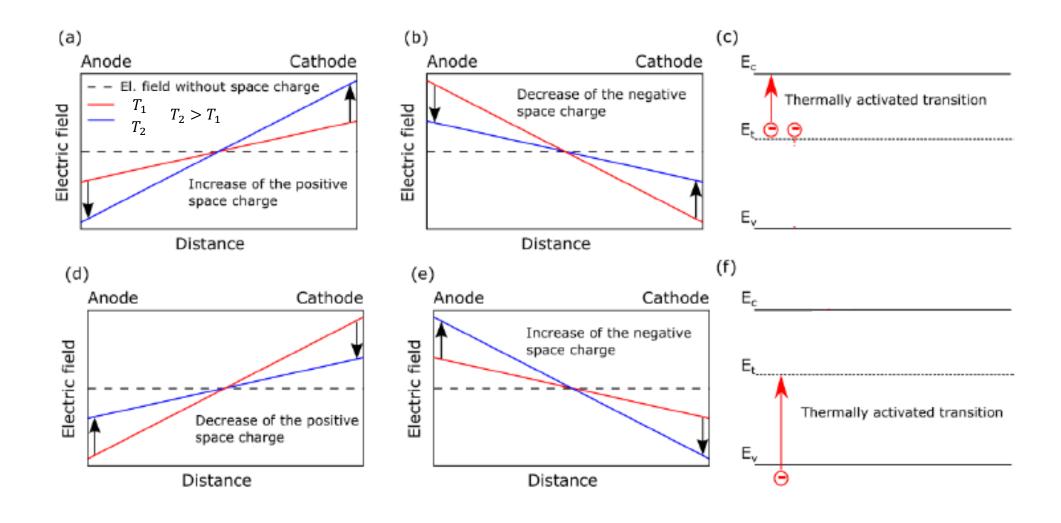
 E_v

$$Q(t) = A \exp\left(-\frac{t}{\tau}\right) + Q^{\infty}.$$

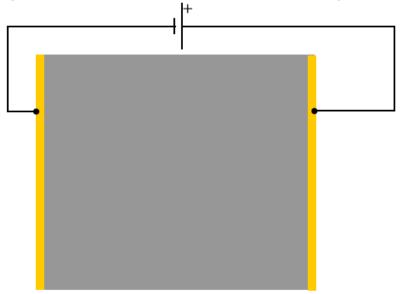


Arrhenius equation:
$$\ln(\tau T^2) = \frac{E}{kT} + \ln\left(\frac{C}{\sigma}\right)$$
 $C = \frac{h^3}{16m_{e(h)}^*\pi k_B^2}$

Determination of type of transitions

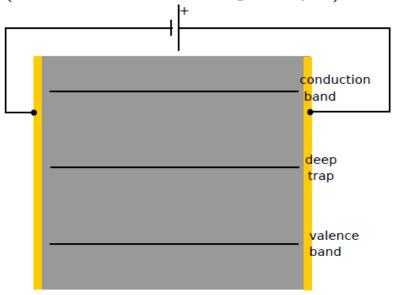


Example (biased semiinsulating sample):



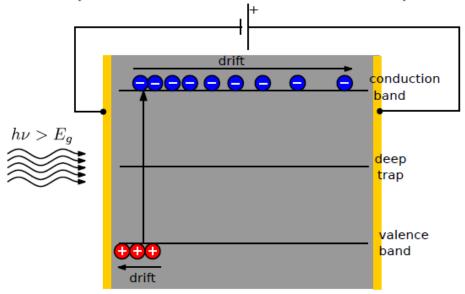
- defect structure (traps, optical transitions)
- penetration depth (photon energy)
- polarity (in case of low penetration)
- electrode material

Example (biased semiinsulating sample):



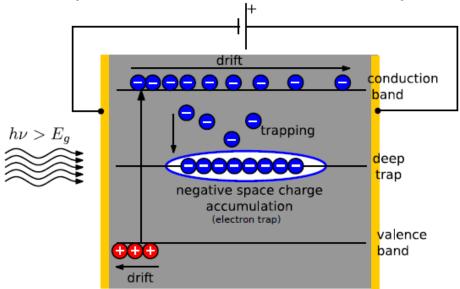
- defect structure (traps, optical transitions)
- penetration depth (photon energy)
- polarity (in case of low penetration)
- electrode material

Example (biased semiinsulating sample):



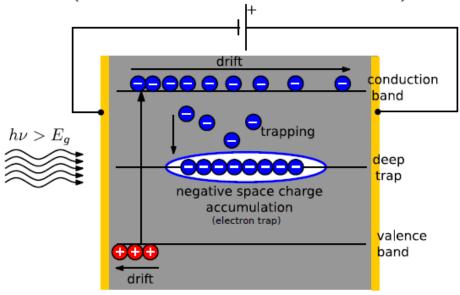
- defect structure (traps, optical transitions)
- penetration depth (photon energy)
- polarity (in case of low penetration)
- electrode material

Example (biased semiinsulating sample):

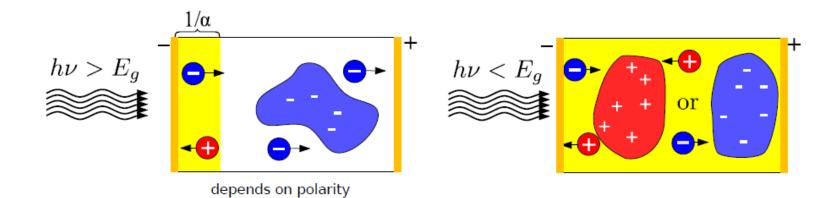


- defect structure (traps, optical transitions)
- penetration depth (photon energy)
- polarity (in case of low penetration)
- electrode material

Example (biased semiinsulating sample):



- defect structure (traps, optical transitions)
- penetration depth (photon energy)
- polarity (in case of low penetration)
- electrode material

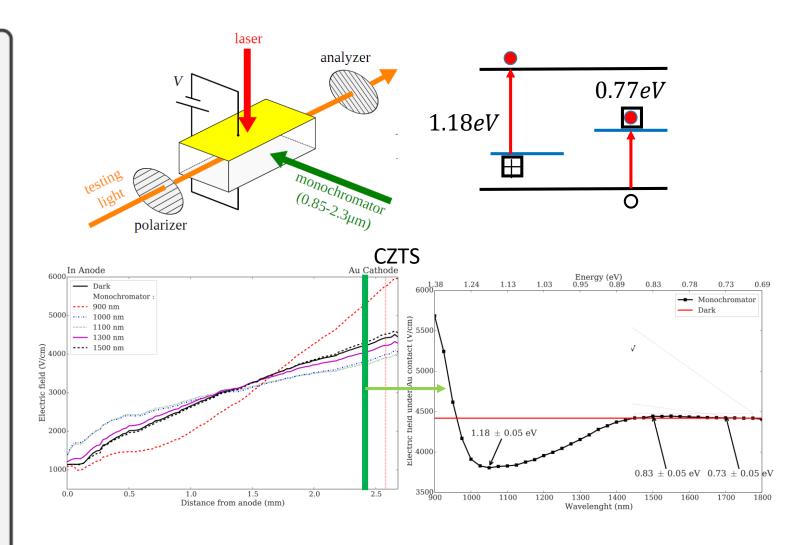


Experimental methods based on Pockels effect II

3-lights experiment

(infrared spectral scanning)

- Subject of study: optical transitions on deep levels
- Light sources:
 - Laser: band-to-band carrier generation (injection)⇒trap filling
 - Tunable monochromatic light: activation of optical transitions
 - Testing light: very low intensity, Pockels effect probe
- Output: E-field (charge) dependence on the photon energy reflecting the transitions enhanced by band-to-band generation (optical injection of carriers)⇒deep levels responsible for space charge accumulation
- Advantage: higher sensitivity than current techniques, direct space charge sign estimation



Electric field profiles X-rays @ 80kVp, 5W LED 1200nm @ 2x1016 cm-2s-1 - X-rays only · · · · X-rays & LED 1200nm preamplifier 3mm thick GND = Al box Monochromator wavelength (nm) d under the cathode (KV/cm) - LED 940 nm only Distance from the cathode (mm) monochromator + LED 940 nm X-rays Z Electric field polarization Monochromator photon energy (eV) X-rays @ 80kVp X-rays @ 80kVp X-ray tube & LED 1200nm @ 50mA Power (W): X-ray tube (flux 2x10¹⁶ cm⁻²s⁻¹) Power (W): polarization Counts x10 pile-ups 120 160 200 120 160 200 240 280 40 80 240 Photon energy (keV) Photon energy (keV)

De-polarization at high fluxes

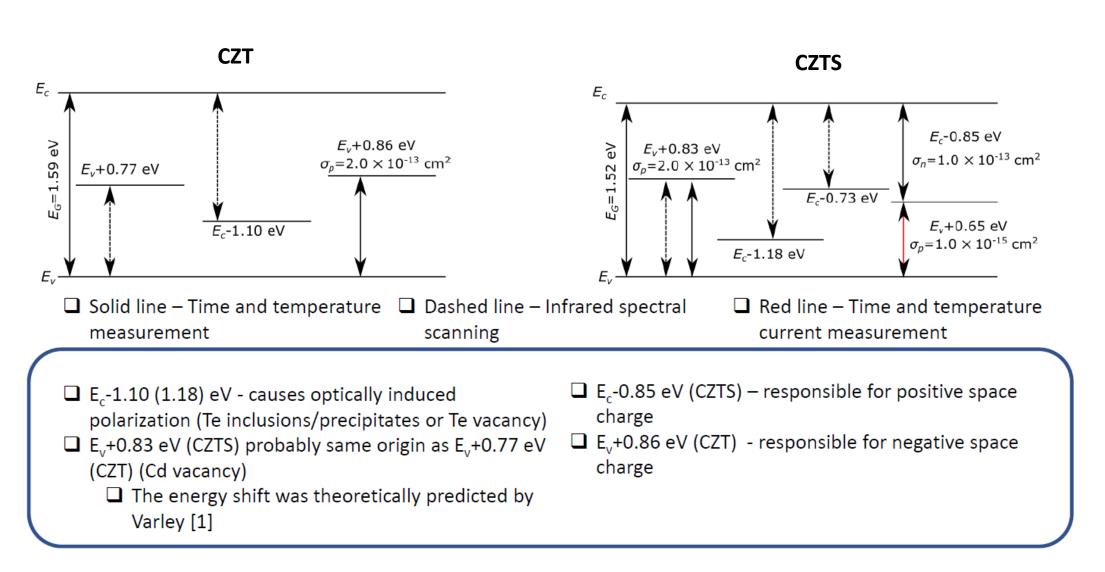
"De-polarization" of high-flux X-ray detector

- X-ray (tube at 80 kV) analogous e-h pair creation to band-to-band generation due to photoeffect
- significant positive space charge accumulation due to strong hole trapping ⇒ screening of the electric field (also "detector polarization")⇒ reduction of CCE
- optical transition reducing space charge⇒ detector recovery



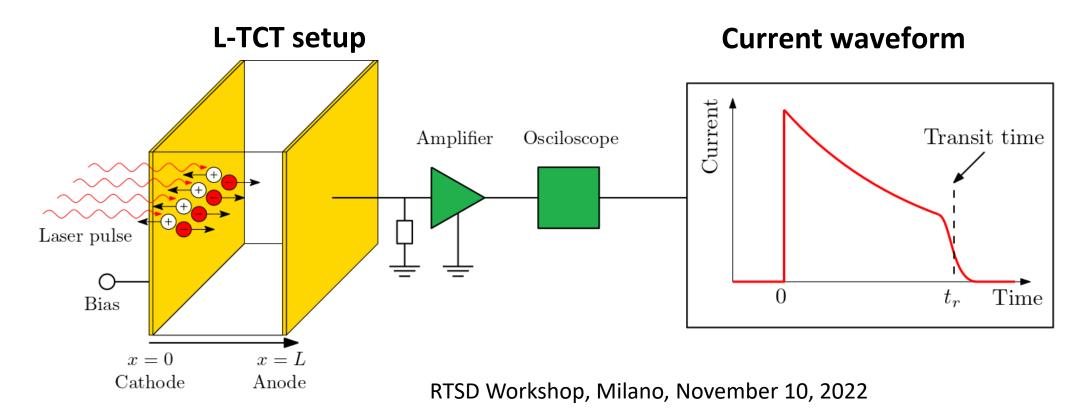
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Summary of deep levels



Laser Induced Current Technique (L-TCT)

- Based on measuring the current response of the detector to a laser pulse
- Allows to characterize the charge transport (mobility, lifetime, electric field profile)
- Bias polarity selects which carrier type drift thought detector



Monte Carlo Simulation

- 1D numerical simulation of charge transport in semiconductor detector
- Combined with numerical solution of driftdiffusion equation and Poisson's equation
- allows study of charge transport and space charge dynamics

Band diagram E_C $\tau_T \quad \tau_D \quad E_T$ $T_T \quad \text{Trapping time}$

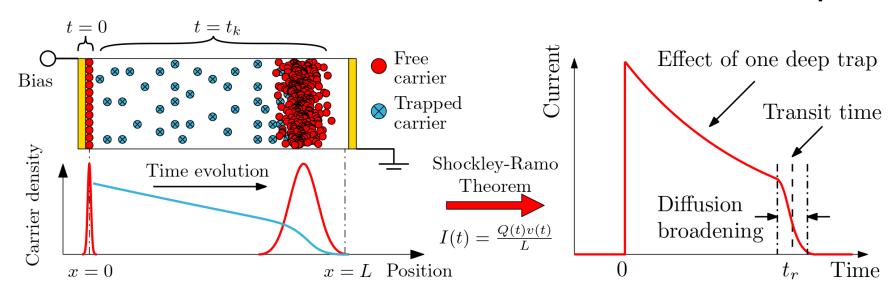
 τ_T Trapping time

 τ_D Detrapping time

 E_V

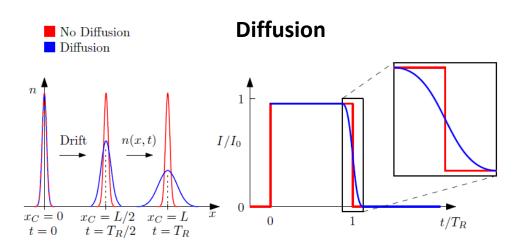
Current response

Detector

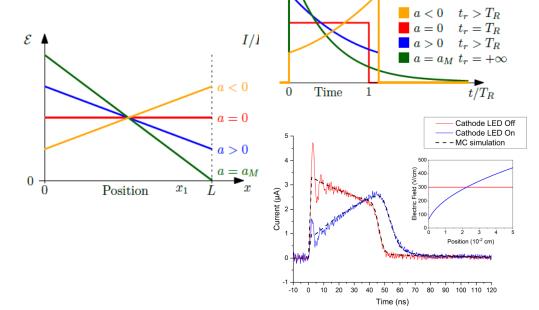


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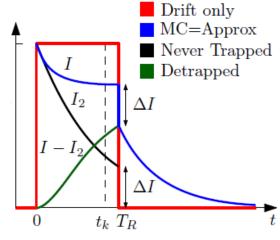
Effects influencing current waveform



Space charge

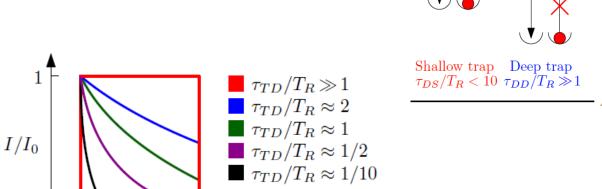


Trapping and de-trapping



$$\tau_T = \frac{1}{N_T \sigma_c v_{th}},$$

$$\tau_D = \frac{1}{N_C \sigma_c v_{th}} \exp\left(\frac{E_T}{k_B T}\right).$$



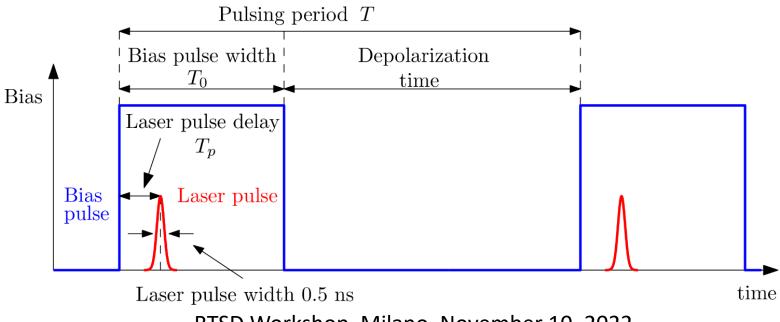
 t/T_R

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Space charge elimination

- Space charge deteriorates detector performance
- Without bias the detector is neutral
- Space charge formation starts after bias application
- Space charge can be eliminated using pulsed bias and effects of space charge can be distinguished from effects of traps
- Changing the laser pulse delay allows study of space charge evolution

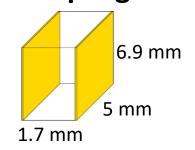
Synchronization of laser pulse and bias pulse



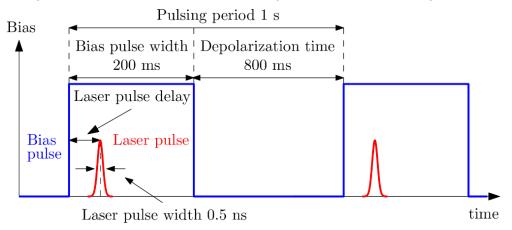
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CdZnTeSe - L-TCT measurement results

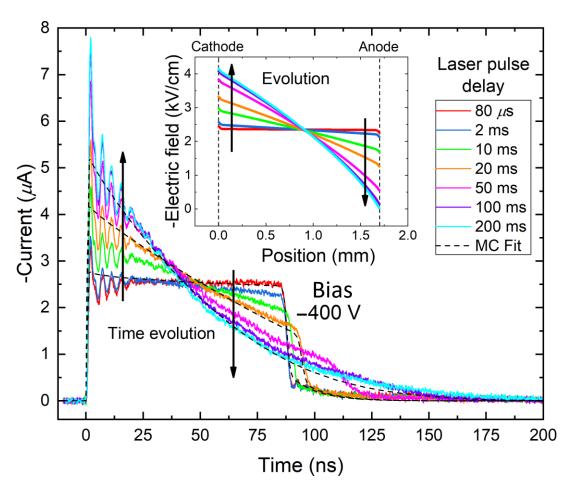
- Semi-insulating p-type Cd_{0.9}Zn_{0.1}Te_{0.96}Se_{0.04} sample is used
- L-TCT is combined with pulsed bias to study space charge dynamics
 Sample geometry



Synchronization of laser pulse and bias pulse



Electron-current waveforms

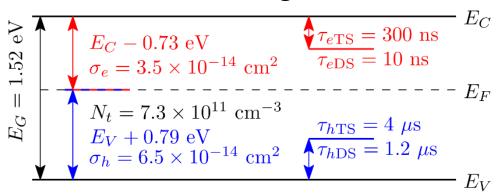


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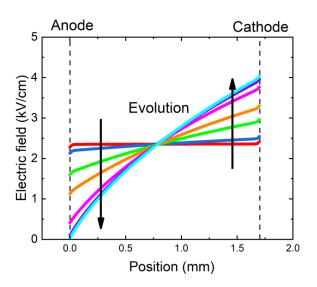
Space charge dynamics

- Three defect levels are sufficient to describe all observed effects
- Positive space charge forms due to hole injection from anode combined with recombination level near Fermi level
- Shallow levels do not contribute to space charge

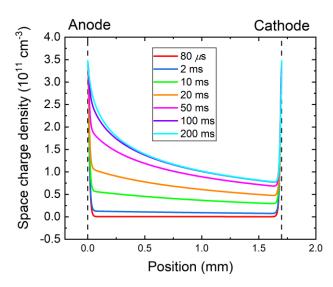
Band diagram



Electric field profile

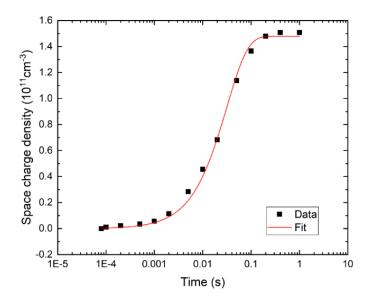


Space charge profile



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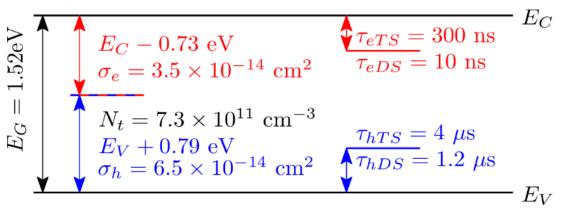
Mean space charge evolution



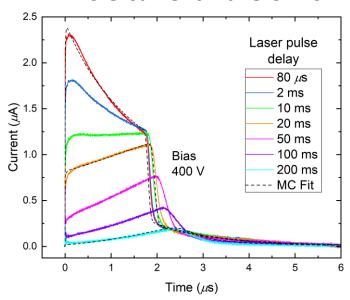
Evaluated parameters

- Same experiment was measured with positive bias for holes
- Identical electric field profile is obtained
- Evaluated parameters are:
- Electron mobility $\mu_e = 830 \text{ cm}^2/\text{Vs}$
- Hole mobility $\mu_h = 40 \text{ cm}^2/\text{Vs}$
- Electron lifetime $\tau_e = 2.3 \,\mu s$
- Hole lifetime $\tau_h = 3.6 \,\mu s$
- Three defect levels:

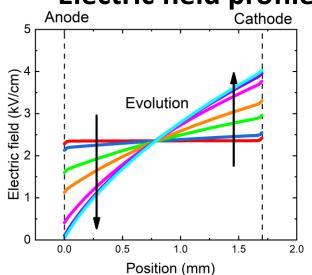
Band diagram



Hole-current waveforms



Electric field profile



Origin of deep levels

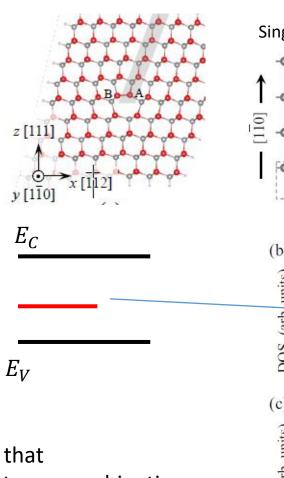
60° Cd core glide dislocation in the (111) plane – first principles calculations

Quadruple

Due to a small concentration of deep levels is an experimental investigation of their origin difficult

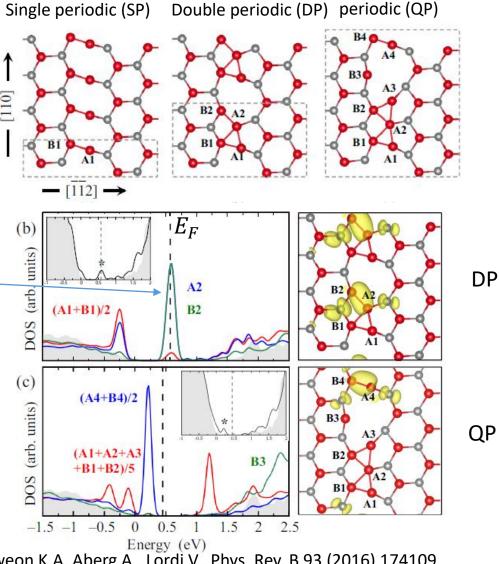
Possible sources:

- Native defects and their complexes (V_{Cd}, Te_{Cd}, ...)
- Impurities
- dislocations



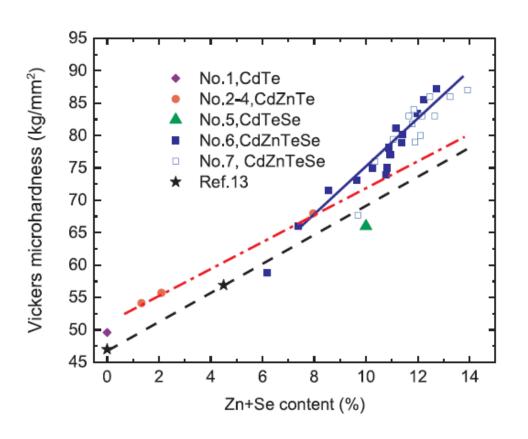
Theoretical calculations of DOS show that dislocation-related deep levels can act as recombination and trapping centers and their position in the gap agree with experiment.

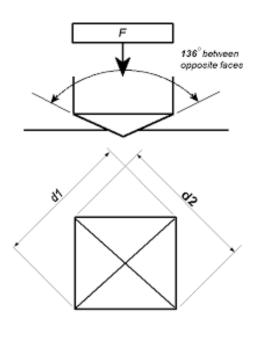
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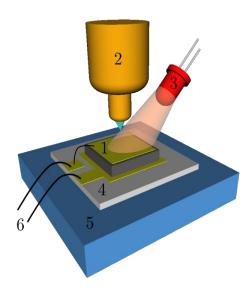
Kweon K.A, Aberg A., Lordi V., Phys. Rev. B 93 (2016) 174109, https://doi.org/10.1103/PhysRevB.93.174109

Vickers microhardness





LED or white laser with bandpass filter

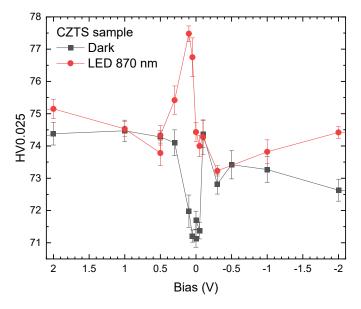


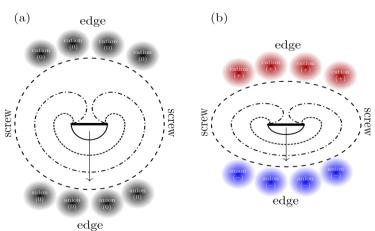
This result shows that Se plays a major role in an effective solution hardening of the CZTS matrix and indicates a possible additional strengthening of the effect when Zn and Se are mixed in the CdTe lattice.

Modified setup – measurement of microhardness with illumination and/or applied bias

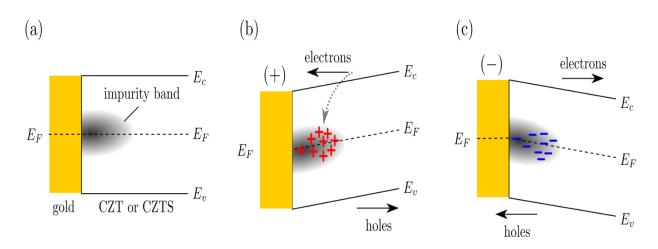
RTSD Workshop, Milano, November 10, 2022

Bias dependence of photo plasticity





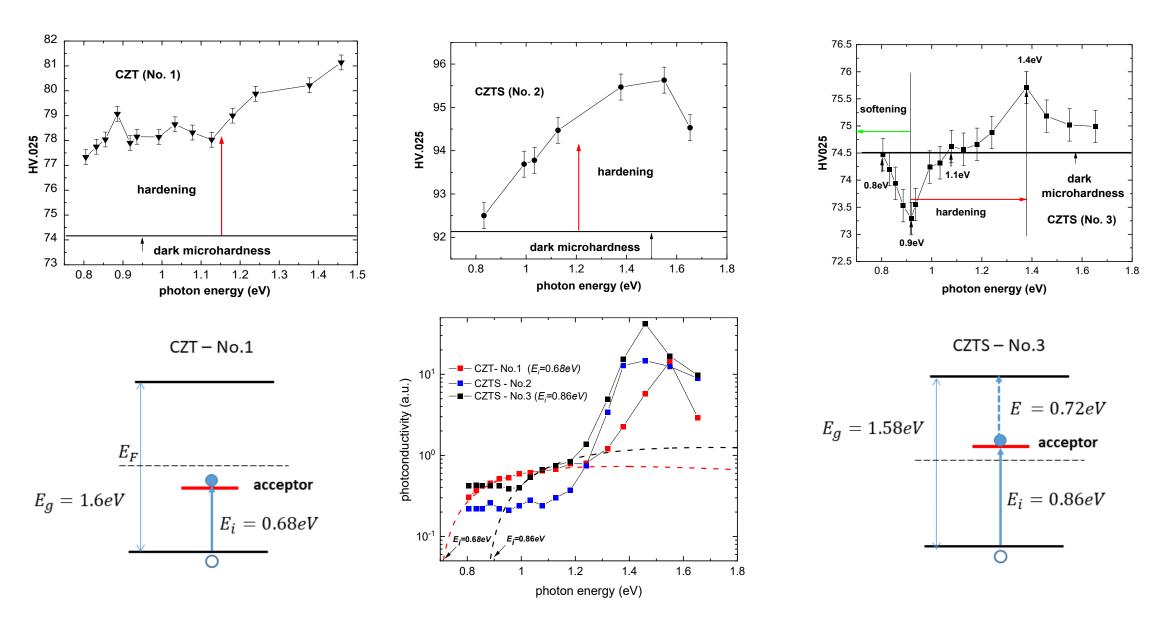
The cation-rich and anion-rich regions accumulate at the opposite sides of the dislocation source. Uncharged dislocations outlined in the left panel (a) glide more easily than charged dislocations shown in the right panel



Illuminating the sample leads to the generation of an excessive amount of free carriers that can be captured by dislocation segments according to their electrical character, resulting in dislocation reconstruction and ensuing hardening of the material

In the case of an applied bias, the deformation-induced defect states are filled by the charges from the adjacent contact, while the opposite charge is not supplied by the bulk.). One of the charge types (positive/negative) starts to dominate the other one. Subsequently, only appropriate dislocation segments may pass the reconstruction. Due to this reason, the hardness stabilizes near the middle of the hardness at 0 V in the dark and for the illumination

Spectral dependence of photo plasticity



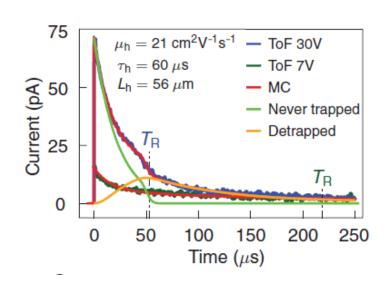
RTSD Workshop, Milano, November 10, 2022

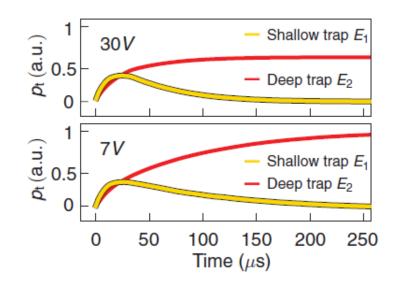
Defects in hybrid perovskites

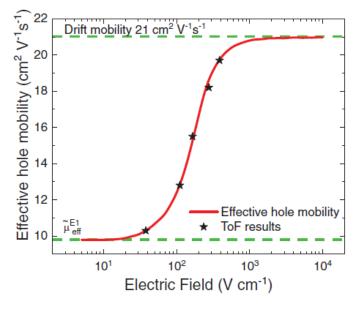
Methylammonium halide perovskites, MAPbI₃, MABrI₃

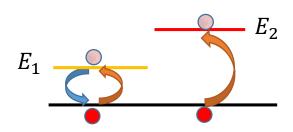
Materials with large application potential – photovoltaics, LED, FET, X-ray and gamma-ray detection **Issues to be solved - stability and defect control**

Hole transport MAPbl₃ – TCT method





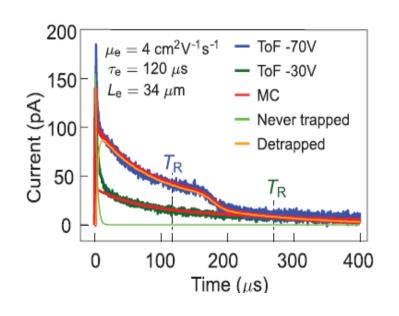


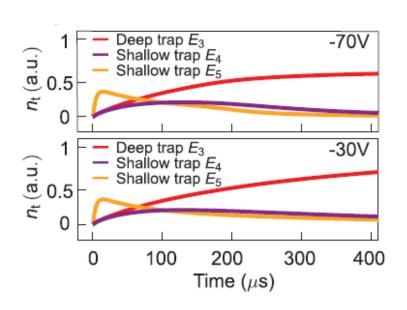


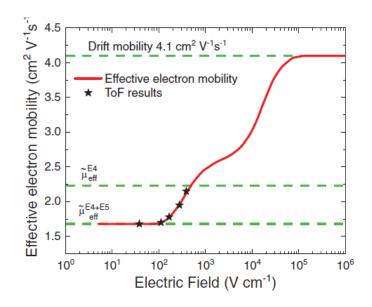
 E_1 is responsible for fast trapping ($\tau_t = 3.5 \mu s$) and de-trapping ($\tau_t = 4 \mu s$) resulting in reduction of effective hole mobility $\mu_{eff} = \frac{\mu_e}{1 + \frac{\tau_D}{\tau_t}}$

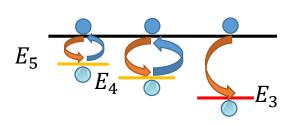
 E_2 permanently traps free holes ($\tau_t=60\mu s$, τ_D » 20ms) resulting in reduction of effective hole mobility

Electron transport MAPbl₃ – **TCT method**







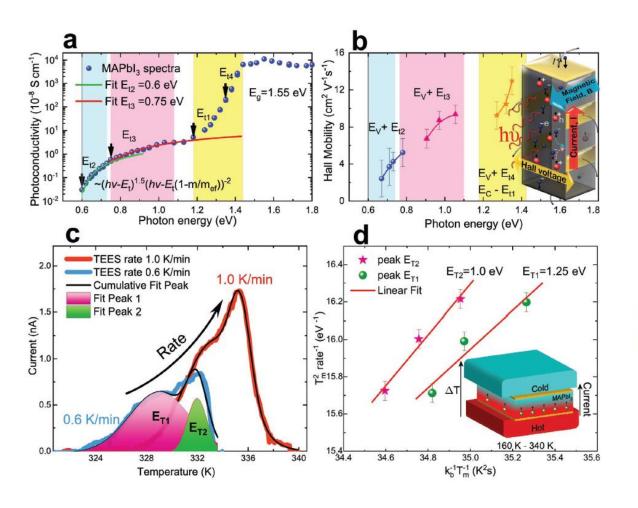


 E_4 and E_5 are responsible for fast trapping

$$\mu_{eff} = \frac{\mu_e}{1 + \frac{\tau_{D4}}{\tau_{t4}} + \frac{\tau_{D5}}{\tau_{t5}}}$$

 E_3 permanently traps free electrons ($\tau_t=120\mu s$, $\tau_D\gg20ms$) resulting in reduction of effective hole mobility

Deep levels in MAPbl₃ – summary

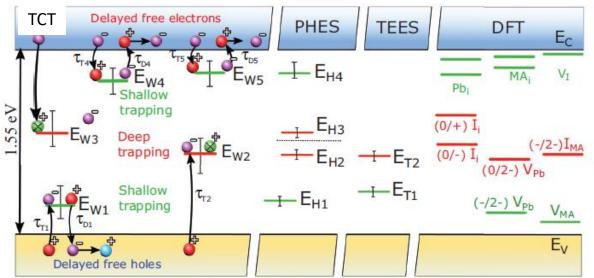


TCT - transient current technique

PHES - photo Hall effect spectroscopy

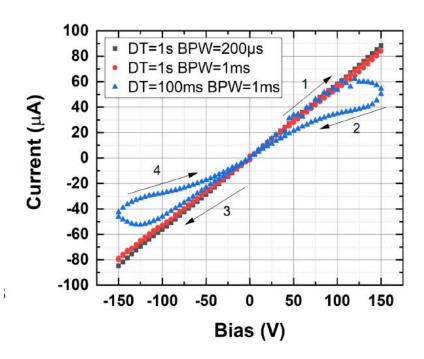
TEES – thermoelectric effect spectroscopy

DFT – density functional theory



A. Musiienko, D.Cerrati, J. Pipek et al, Adv. Func. Mat. 2021, 31, 2104467

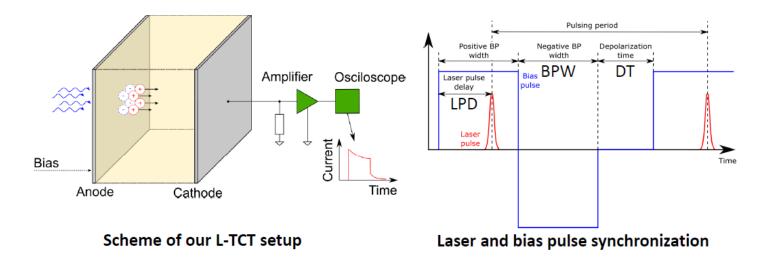
Space charge formation nad relaxation in MAPbBr₃

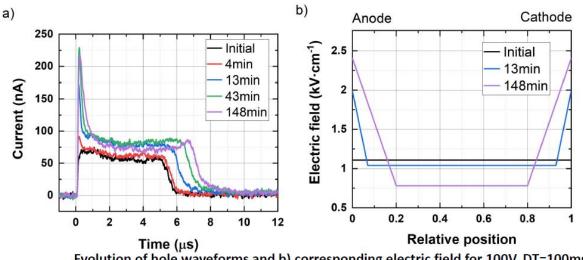


Polarization effects observed

Two models to explain:

- 1.lon migration
- 2. Trapping of carriesr at deep levels

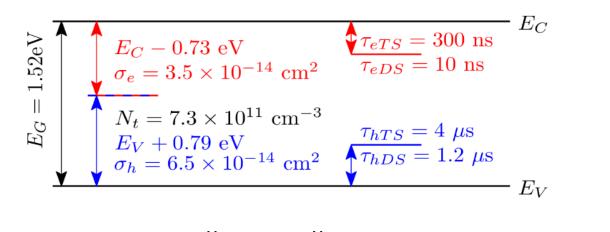




Evolution of hole waveforms and b) corresponding electric field for 100V, DT=100ms BPW=1ms and LPD=100µs.

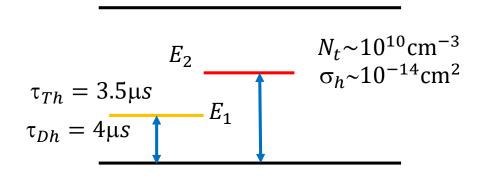
Level comparison

CdZnTeSe group



$$\mu_{eff,e} = \frac{\mu_e}{1 + \frac{\tau_D}{\tau_t}} = \frac{\mu_e}{1 + \frac{10}{300}} = 0.967\mu_e$$





$$\mu_{eff,h} = \frac{\mu_e}{1 + \frac{\tau_D}{\tau_t}} = \frac{\mu_e}{1 + \frac{4}{3.5}} = 0.467\mu_e$$

Midgap level - similar concentration, trapping activity potentially leading to polarization

Shallow traps – in both cases the effective mobility is decreased by trapping at de-trapping at deep levels. The effect is stronger in measured perovskites samples

Future work

Development of L-TCT setup (temperature dependence measurements) and its applications to perovskites, CdTe group compounds, and other materials

Development of Pockels effect method - measurements of graphene/SiC structures, perovskites etc.

Continue to study mechanical properties – Vickers indentation, AFM - CdTe group, perovskites

Thank you very much for your attention